

Finite Difference Method

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September 26, 2015

Introduction

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Math Modeling and Simulation of Physical Processes

- ▶ Describe the physical phenomenon
- ▶ Model the physical phenomenon to become mathematical equations(PDE)
- ▶ Simulate the mathematic equations (discrete solution)
- ▶ Compare the discrete solution and experiment result

Some kind of Partial Differential Equation (PDE)

- ▶ Elliptic equation
 - ▶ Diffusion equation
 - ▶ Poisson's equation
- ▶ Parabolic equation
 - ▶ Heat equations
- ▶ Hyperbolic equation
 - ▶ Wave equation
 - ▶ The equation for conservation laws

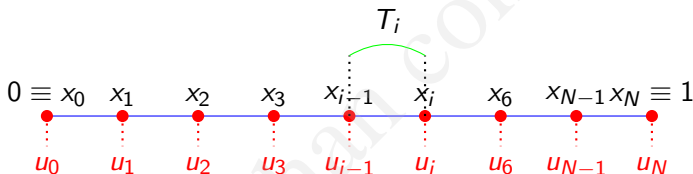
Laplace equation

We consider the partial differential equation on $]0, 1[$

$$\begin{cases} -u_{xx}(x) &= f(x) & \text{for all } x \in]0, 1[\\ u(0) &= 0 \\ u(1) &= 0 \end{cases} \quad (1)$$

To find the discrete solution of this equation, there are many methods, we will choose a method which is the simplest method, it is the finite difference scheme.

Mesh



Let us consider a uniform partition with $N + 1$ points x_i for all $i = 0, 1, 2, \dots, N$ (see figure). We have space step is $\Delta x = \frac{1}{N}$, then

$$x_i = i\Delta x$$

Our purpose is the value of the function at points x_i

$$u_i \simeq u(x_i) \text{ for all } i = 0, 1, 2, \dots, N$$

Approximation of derivatives

$$\frac{\partial u}{\partial x}(x_i) = \frac{u_{i+1} - u_i}{\Delta x} \text{ forward difference}$$

$$\frac{\partial u}{\partial x}(x_i) = \frac{u_i - u_{i-1}}{\Delta x} \text{ backward difference}$$

$$\frac{\partial u}{\partial x}(x_i) = \frac{u_{i+1} - u_{i-1}}{2\Delta x} \text{ central difference}$$

Approximation of derivatives (Cont.)

Use the Taylor series expansion at x_i

$$\begin{aligned}u(x_{i+1}) = & u(x_i) + \frac{\partial u}{\partial x}(x_i)(x_{i+1} - x_i) + \frac{\frac{\partial^2 u}{\partial x^2}(x_i)}{2!}(x_{i+1} - x_i)^2 \\& + \frac{\frac{\partial^3 u}{\partial x^3}(x_i)}{3!}(x_{i+1} - x_i)^3 + O((x_{i+1} - x_i)^4)\end{aligned}$$

Or

$$u_{i+1} = u_i + \frac{\partial u}{\partial x}(x_i)\Delta x + \frac{\frac{\partial^2 u}{\partial x^2}(x_i)}{2!}\Delta^2 x + \frac{\frac{\partial^3 u}{\partial x^3}(x_i)}{3!}\Delta^3 x + O(\Delta^4 x) \quad (2)$$

We can approximate the derivative $\frac{\partial u}{\partial x}(x_i)$ that

$$\frac{\partial u}{\partial x}(x_i) = \frac{u_{i+1} - u_i}{\Delta x} + O(\Delta x)$$

Approximation of derivatives

It is similar, we obtain

$$u_{i-1} = u_i - \frac{\partial u}{\partial x}(x_i)\Delta x + \frac{\frac{\partial^2 u}{\partial x^2}(x_i)}{2!}\Delta^2 x - \frac{\frac{\partial^3 u}{\partial x^3}(x_i)}{3!}\Delta^3 x + 0(\Delta^4 x) \quad (3)$$

We can approximate the derivative $\frac{\partial u}{\partial x}(x_i)$ that

$$\frac{\partial u}{\partial x}(x_i) = \frac{u_i - u_{i-1}}{\Delta x} + 0(\Delta x)$$

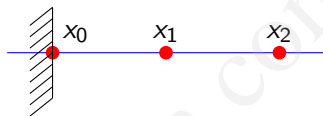
Let (2)-(3), we have

$$u_{i+1} - u_{i-1} = 2\frac{\partial u}{\partial x}(x_i)\Delta x + 2\frac{\frac{\partial^3 u}{\partial x^3}(x_i)}{3!}\Delta^3 x + 0(\Delta^4 x)$$

We can also approximate the derivative $\frac{\partial u}{\partial x}(x_i)$ that

$$\frac{\partial u}{\partial x}(x_i) = \frac{u_{i+1} - u_{i-1}}{2\Delta x} + 0(\Delta^2 x)$$

Approximation of derivative at boundary



We use the Taylor series expansion at x_0

$$u(x_1) = u(x_0) + \frac{\partial u}{\partial x}(x_0)(x_1 - x_0) + \frac{\partial^2 u}{\partial x^2}(x_0) \frac{(x_1 - x_0)^2}{2!} + O((x_1 - x_0)^3)$$

Or

$$u(x_1) = u(x_0) + \frac{\partial u}{\partial x}(x_0)\Delta x + \frac{\partial^2 u}{\partial x^2}(x_0) \frac{\Delta x^2}{2!} + O(\Delta x^3) \quad (4)$$

And

$$u(x_2) = u(x_0) + 2\frac{\partial u}{\partial x}(x_0)\Delta x + 2\frac{\partial^2 u}{\partial x^2}(x_0) \frac{\Delta x^2}{2!} + O(\Delta x^3) \quad (5)$$

Approximation of the derivatives at boundary (Cont.)

From (4), we have

$$\begin{aligned}\frac{\partial u}{\partial x}(x_0) &= \frac{u(x_1) - u(x_0)}{\Delta x} + O(\Delta x) \\ &= \frac{u_1 - u_0}{\Delta x}\end{aligned}\quad (6)$$

Combining (4) and (5), there holds

$$u(x_2) - 4u(x_1) = -3u(x_0) - 2\frac{\partial u}{\partial x}(x_0) + O(\Delta^3 x)$$

or

$$\frac{\partial u}{\partial x}(x_0) = \frac{-3u_0 + 4u_1 - u_2}{2\Delta x} + O(\Delta^2 x) \quad (7)$$

Approximation of the second order derivatives

Using again the Taylor series expansion, there holds

$$u_{i+1} = u_i + \frac{\partial u}{\partial x}(x_i)\Delta x + \frac{\frac{\partial^2 u}{\partial x^2}(x_i)}{2!}\Delta^2 x + \frac{\frac{\partial^3 u}{\partial x^3}(x_i)}{3!}\Delta^3 x + o(\Delta^4 x)$$

and

$$u_{i-1} = u_i - \frac{\partial u}{\partial x}(x_i)\Delta x + \frac{\frac{\partial^2 u}{\partial x^2}(x_i)}{2!}\Delta^2 x - \frac{\frac{\partial^3 u}{\partial x^3}(x_i)}{3!}\Delta^3 x + o(\Delta^4 x)$$

Adding two previous approximate equations side by side, we have

$$u_{i+1} + u_{i-1} = 2u_i + \frac{\partial^2 u}{\partial x^2}(x_i)\Delta^2 x + o(\Delta^4 x) \quad (8)$$

or

$$\frac{\partial^2 u}{\partial x^2}(x_i) = \frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta^2 x} + o(\Delta^2 x) \quad (9)$$

Discretizing Laplace equation

From the first equation of (1), we have

$$-\frac{\partial^2 u}{\partial x^2}(x_i) = f(x_i) \quad \text{for all } i = 1, \dots, N-1$$

Using the approximation in (9), there holds

$$-\frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta^2 x} = f_i \quad \text{for all } i = 1, \dots, N-1, \quad (10)$$

where $f_i = f(x_i)$ for $i = 1, \dots, N-1$.

Using the Dirichlet boundary condition, we obtain

$$u_0 = 0 \quad \text{and} \quad u_N = 0$$

Discrete equations

Linear system for the scheme

$$\left\{ \begin{array}{ll} i = 1, \frac{2u_1 - u_2}{\Delta^2 x} & = f_1 \\ i = 2, \frac{-u_1 + 2u_2 - u_3}{\Delta^2 x} & = f_2 \\ i = 3, \frac{-u_2 + 2u_3 - u_4}{\Delta^2 x} & = f_3 \\ \vdots & \\ i = N-2, \frac{-u_{N-2} + 2u_{N-1} - u_N}{\Delta^2 x} & = f_{N-2} \\ i = N-1, \frac{-u_{N-1} + 2u_N - u_{N+1}}{\Delta^2 x} & = f_{N-1} \end{array} \right.$$

Matrix form $AU = F$, $A \in \mathbb{R}^N \times \mathbb{R}^N$, $U, F \in \mathbb{R}^N$,

$$A = \frac{1}{\Delta^2 x} \begin{bmatrix} 2 & -1 & 0 & 0 & 0 & 0 \\ -1 & 2 & -1 & 0 & 0 & 0 \\ 0 & -1 & 2 & -1 & 0 & 0 \\ 0 & 0 & 0 & \dots & 0 & 0 \\ 0 & 0 & 0 & -1 & 2 & -1 \\ 0 & 0 & 0 & 0 & -1 & 2 \end{bmatrix}$$

$$U = \begin{bmatrix} u_1 \\ u_2 \\ u_3 \\ \vdots \\ u_{N-2} \\ u_{N-1} \end{bmatrix} \quad F = \begin{bmatrix} f_1 \\ f_2 \\ f_3 \\ \vdots \\ f_{N-2} \\ f_{N-1} \end{bmatrix}$$

The matrix A remains tridiagonal and symmetric positive definite

Other types of boundary condition

■ Dirichlet Neumann Boundary Condition: $u(0) = \frac{\partial u}{\partial x}(1) = 0$.

► Using the backward difference at 1, it means that

$$\frac{\partial u}{\partial x}(1) = \frac{u_N - u_{N-1}}{\Delta x} = 0 \Rightarrow u_{N-1} = u_N$$

Only changing the last equation in the linear system:

$$\frac{-u_{N-2} + u_{N-1}}{\Delta^2 x} = f_{N-1}$$

Other types of boundary condition

Then the linear system for the scheme

$$\left\{ \begin{array}{ll} i = 1, \frac{2u_1 - u_2}{\Delta^2 x} & = f_1 \\ i = 2, \frac{-u_1 + 2u_2 - u_3}{\Delta^2 x} & = f_2 \\ i = 3, \frac{-u_2 + 2u_3 - u_4}{\Delta^2 x} & = f_3 \\ & \dots \\ i = N - 2, \frac{-u_{N-3} + 2u_{N-2} - u_{N-1}}{\Delta^2 x} & = f_{N-2} \\ i = N - 1, \frac{-u_{N-2} + u_{N-1}}{\Delta^2 x} & = f_{N-1} \end{array} \right.$$

Other types of boundary condition (Cont.)

- Using the second order approximation of the derivative at 1, it means that

$$\frac{\partial u}{\partial x}(1) = \frac{-3u_N + 4u_{N-1} - u_{N-2}}{2\Delta x} = 0$$

Implied

$$u_N = \frac{4u_{N-1} - u_{N-2}}{3}$$

Changing only the last equation in the linear system, the last equation becomes

$$\frac{-u_{N-2} + u_{N-1}}{\Delta^2 x} = \frac{3}{2}f_{N-1}$$

Other types of boundary condition (Cont.)

Then the linear system for the scheme

$$\left\{ \begin{array}{ll} i = 1, \frac{2u_1 - u_2}{\Delta^2 x} & = f_1 \\ i = 2, \frac{-u_1 + 2u_2 - u_3}{\Delta^2 x} & = f_2 \\ i = 3, \frac{-u_2 + 2u_3 - u_4}{\Delta^2 x} & = f_3 \\ & \dots \\ i = N-2, \frac{-u_{N-3} + 2u_{N-2} - u_{N-1}}{\Delta^2 x} & = f_{N-2} \\ i = N-1, \frac{-u_{N-2} + u_{N-1}}{\Delta^2 x} & = \frac{3}{2} f_{N-1} \end{array} \right.$$

Other types of boundary condition (Cont.)

- ▶ Using the central difference at 1, it means that

$$\frac{\partial u}{\partial x}(1) = \frac{u_{N+1} - u_{N-1}}{2\Delta x}$$

Implied

$$u_{N+1} = u_{N-1}$$

We discretize additionally at point $x_N = 1$, there holds

$$\frac{-u_{N-1} + 2u_N - u_{N+1}}{\Delta^2 x} = f_N$$

where $f_N = f(x_N)$. Combining with discrete boundary condition, we have

$$\frac{-u_{N-1} + u_N}{\Delta^2 x} = \frac{f_N}{2}$$

Other types of boundary condition (Cont.)

Then the linear system for the scheme

$$\left\{ \begin{array}{ll} i = 1, \frac{2u_1 - u_2}{\Delta^2 x} & = f_1 \\ i = 2, \frac{-u_1 + 2u_2 - u_3}{\Delta^2 x} & = f_2 \\ i = 3, \frac{-u_2 + 2u_3 - u_4}{\Delta^2 x} & = f_3 \\ & \dots \\ i = N-1, \frac{-u_{N-2} + 2u_{N-1} - u_N}{\Delta^2 x} & = f_{N-1} \\ i = N, \frac{-u_{N-1} + u_N}{\Delta^2 x} & = \frac{1}{2}f_N \end{array} \right.$$

Other types of boundary condition (Cont.)

■ Non-homogeneous Dirichlet Boundary Condition:

$$u(0) = \alpha, \quad u(1) = \beta.$$

The first and last equations will be changed in the linear system, it means that

$$u_0 = \alpha \Rightarrow \frac{2u_1 - u_2}{\Delta^2 x} = f_1 + \frac{\alpha}{\Delta^2 x},$$
$$u_N = \beta \Rightarrow \frac{-u_{N-2} + 2u_{N-1}}{\Delta^2 x} = f_{N-1} + \frac{\beta}{\Delta^2 x}$$

Other types of boundary condition (Cont.)

Then the linear system for the scheme

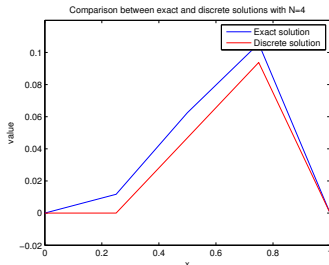
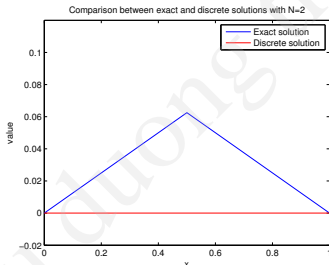
$$\left\{ \begin{array}{ll} i = 1, \frac{2u_1 - u_2}{\Delta^2 x} & = f_1 + \frac{\alpha}{\Delta^2 x} \\ i = 2, \frac{-u_1 + 2u_2 - u_3}{\Delta^2 x} & = f_2 \\ i = 3, \frac{-u_2 + 2u_3 - u_4}{\Delta^2 x} & = f_3 \\ & \dots \\ i = N-2, \frac{-u_{N-3} + 2u_{N-2} - u_{N-1}}{\Delta^2 x} & = f_{N-2} \\ i = N-1, \frac{-u_{N-2} + 2u_{N-1}}{\Delta^2 x} & = f_{N-1} + \frac{\beta}{\Delta^2 x} \end{array} \right.$$

Experiment test

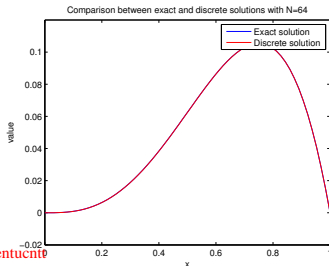
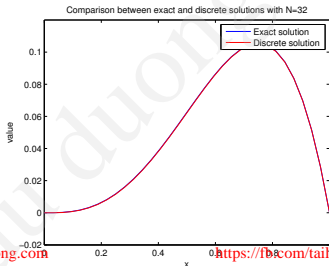
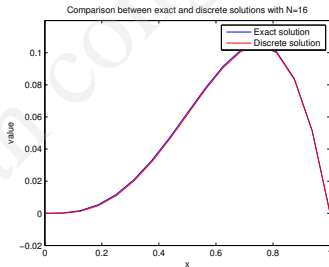
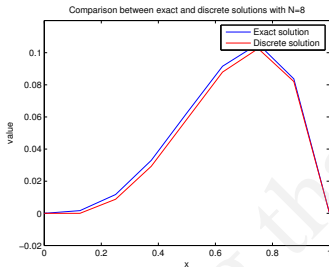
We set up with the following exact solution $u(x)$ and function $f(x)$

$$f(x) = 12x^2 - 6x$$

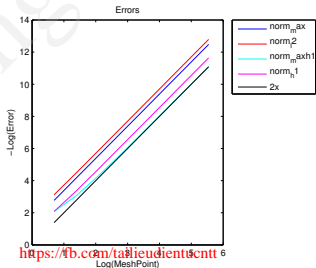
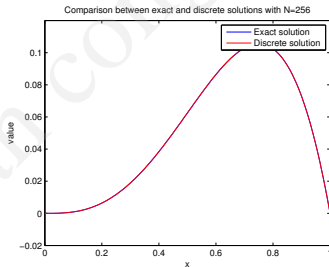
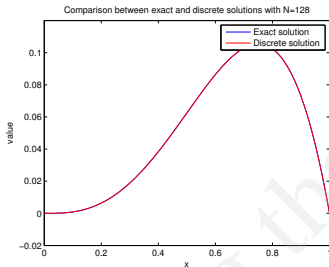
$$u(x) = x^3(1 - x)$$



Experiment test



Experiment test



Norms

We definite

$$U = \begin{bmatrix} u_0 \\ u_1 \\ u_2 \\ \vdots \\ u_{N-1} \\ u_N \end{bmatrix} \quad \text{and} \quad \hat{U} = \begin{bmatrix} u(x_0) \\ u(x_1) \\ u(x_1) \\ \vdots \\ u(x_{N-1}) \\ u(x_N) \end{bmatrix}$$

and Error $E = U - \hat{U}$ containt the errors at each grid point.

To estimate the amplitude of error vector, we define somes norm on it.

Definition (L_h^∞ -norm)

$$\|E\|_{\infty, h} = \max_{0 \leq i \leq N} |E_i| = \max_{0 \leq i \leq N} |u_i - u(x_i)|$$

Norms

We put $h_i = |x_{i+1} - x_i|$ for all $i = 0, \dots, N-1$

Definition (L_h^1 -norm)

$$\|E\|_{1,h}^+ = \sum_{i=0}^{N-1} |E_i| h_i = \sum_{i=0}^N |u_i - u(x_i)| h_i$$

$$\|E\|_{1,h}^- = \sum_{i=1}^N |E_i| h_{i-1} = \sum_{i=1}^N |u_i - u(x_i)| h_{i-1}$$

Definition (L_h^2 -norm)

$$\|E\|_{2,h}^+ = \sum_{i=0}^{N-1} |E_i|^2 h_i = \sum_{i=1}^N |u_i - u(x_i)|^2 h_i$$

$$\|E\|_{2,h}^- = \sum_{i=1}^N |E_i|^2 h_{i-1} = \sum_{i=1}^N |u_i - u(x_i)|^2 h_{i-1}$$

Local Truncation Error

We can replace discrete solution u_i by exact solution $u(x_i)$ in (10). In general, the exact solution won't satisfy this equation, which define τ_i

$$\tau_i = -\frac{1}{h^2}(u(x_{i-1}) - 2u(x_i) + u(x_{i+1})) - f(x_i) \text{ for all } i = 1, \dots, N-1 \quad (11)$$

Using Taylor series, we get

$$\tau_i = -\left[u''(x_i) + \frac{1}{12}h^2 u''''(x_i) + O(h^4)\right] - f(x_i) \quad (12)$$

Using our original differential equation (1) this becomes

$$\tau_i = -\frac{1}{12}h^2 u''''(x_i) - O(h^4) = O(h^2)$$

Global Error

We define τ to be the vector with component τ_i then

$$\tau = A\hat{U} - F \quad (13)$$

also

$$A\hat{U} = \tau + F \quad (14)$$

To obtain a relation between the local error τ and the global error $E = U - \hat{U}$, we get

$$AE = -\tau \quad (15)$$

This is simply the matrix form of the system of equations

$$\frac{1}{h^2}(E_{i-1} - 2E_i + E_{i+1}) = -\tau_i \text{ for all } i \in [1, N-1] \quad (16)$$

with the boundary conditions

$$E_0 = E_N = 0 \quad (17)$$

Let A^{-1} be the inverse of the matrix A . Then solving the system (15) gives

$$E = -A^{-1}\tau$$

and taking norms gives

$$\|E\| = \|A^{-1}\tau\| \leq \|A^{-1}\| \|\tau\| \quad (18)$$

We know that $\|\tau\| = O(h^2)$ and we are hoping the same will be true of $\|E\| = O(h^2)$. It is clear what we need for this to be true: we need $\|A^{-1}\|$ to be bounded by some constant independent of h as $h \rightarrow 0$:

$$\|A^{-1}\| \leq C \text{ for } h \text{ sufficiently small}$$

Stability

Then we will have

$$\|E\| \leq C\|\tau\| \quad (19)$$

so $\|E\|$ goes to zero at least as fast as $\|\tau\|$.

Definition

Suppose a finite difference method for Laplace equation gives a sequence of matrix equations of the form $AU = F$. We say that the method is stable if A^{-1} exists for all h sufficiently small (for $h < h_0$, say) and if there is a constant C , independent of h , such that

$$\|A^{-1}\| \leq C \text{ for all } h < h_0 \quad (20)$$

Consistency

We say that a method is consistent with the differential equation and boundary conditions if

$$\|\tau\| \rightarrow 0 \text{ as } h \rightarrow 0 \quad (21)$$

Convergence

A method is said to be convergent if $\|E\| \rightarrow 0$ as $h \rightarrow 0$.

Combining the ideas introduced above we arrive at the conclusion that

$$\text{consistency} + \text{stability} \implies \text{convergence} \quad (22)$$

This is easily proved by using (20) and (21) to obtain the bound

$$\|E\| \leq \|A^{-1}\| \|\tau\| \leq C \|\tau\| \rightarrow 0 \text{ as } h \rightarrow 0 \quad (23)$$

Stability in L^2 norm

Since the matrix A is symmetric, the L_h^2 -norm of A is equal to its spectral radius

$$\|A\|_{2,h} = \rho(A) = \max_{1 \leq p \leq N-1} \lambda_p \quad (24)$$

where λ_p refers to the p th eigenvalue of the matrix A .

The matrix A^{-1} is also symmetric, and the eigenvalues of A^{-1} are simply the inverses of the eigenvalues of A , so

$$\|A^{-1}\|_{2,h} = \max_{1 \leq p \leq N-1} \lambda_p^{-1} = \left(\min_{1 \leq p \leq N-1} \lambda_p \right)^{-1} \quad (25)$$

So all we need to do is compute the eigenvalues of A and show that they are bounded away from zero as $h \rightarrow 0$

Stability in L^2 norm

We will now focus on one particular value of $h = \frac{1}{N}$. Then the $N - 1$ eigenvalues of A are given by

$$\lambda_p = \frac{2}{h^2}(1 - \cos(\pi ph)) \text{ for all } p = 1, \dots, N - 1 \quad (26)$$

The eigenvector u^p corresponding to p has components u_j^p for $j = 1, \dots, N - 1$ given by

$$u_j^p = \sin(\pi pjh) \quad (27)$$

This can be verified by checking that $Au^p = \lambda_p u^p$. The j th component of the vector Au^p is

Stability in L^2 norm

$$\begin{aligned}(Au^p)_j &= -\frac{1}{h^2}(u_{j-1}^p - 2u_j^p + u_{j+1}^p) \\&= -\frac{1}{h^2}(\sin(\pi p(j-1)h) - 2\sin(\pi pjh) + \sin(\pi p(j+1)h)) \\&= -\frac{1}{h^2}(2\sin(\pi pjh)\cos(\pi ph) - 2\sin(\pi pjh)) \\&= \lambda_p u_j^p\end{aligned}$$

From (26), we see that the smallest eigenvalue of A is

$$\begin{aligned}\lambda_1 &= \frac{2}{h^2}(1 - \cos(\pi h)) \\&= \frac{2}{h^2}\left(\frac{1}{2}\pi^2 h^2 - \frac{1}{24}\pi^4 h^4 + O(h^6)\right) \\&= \pi^2 + O(h^2)\end{aligned}$$

Stability in L^2 norm

This is clearly bounded away from zero as $h \rightarrow 0$, so we see that the method is stable in the L_h^2 -norm. Moreover we get an error bound from this:

$$\|E\|_{2,h} \leq \|A^{-1}\|_{2,h} \|\tau\|_{2,h} \approx \frac{1}{\pi^2} \|\tau\|_{2,h} \quad (28)$$

Since $\tau_j \approx \frac{h^2}{12} u''''(x_j)$, we expect $\|\tau\|_{2,h} \approx \frac{h^2}{12} \|u''''\|_{2,h} = \frac{h^2}{12} \|f''\|_{2,h}$

Stability

we define discrete L_h^2 -norm

$$\|u\|_{2,h}^2 = \sum_{i=0}^{N-1} u_i^2 h$$

Multiplying (10) by u_i then sum over $i = \dots, N-1$, we get

$$\sum_{i=1}^{N-1} \frac{(u_i - u_{i-1})u_i}{h^2} + \frac{(u_i - u_{i+1,j})u_{i,j}}{h^2} = \sum_{i=1}^{N-1} f_i u_i$$

We can change the index in the sum, we have

$$\sum_{i=1}^{N-1} \frac{(u_i - u_{i-1})u_i}{h^2} + \sum_{i=2}^N \frac{(u_{i-1} - u_i)u_{i-1}}{h^2} = \sum_{i=1} f_i u_i$$

Stability

Sine $u_0 = u_N = 0$, then

$$\sum_{i=1}^N \frac{(u_i - u_{i-1})^2}{h^2} = \sum_i^{N-1} f_i u_i$$

We can write again

$$\sum_{i=1}^N (D_{x-} u)_i^2 = \sum_{i=1}^{N-1} f_i u_i, \quad (29)$$

where

$$(D_{x-} u)_i = \frac{u_i - u_{i-1}}{h}$$

Let's define the discrete H_h^1 -norm

$$|||u|||_{1,h}^2 = \sum_{i=1}^N (D_{x-} u)_{i,j}^2 h$$

Stability

Applying Holder inequality, there hold

$$h \sum_{i=1}^{N-1} f_i u_i \leq \left(\sum_{i=0}^{N-1} h f_i^2 \right)^{1/2} \left(\sum_{i=0}^{N-1} h u_i^2 \right)^{1/2} = \|f\|_{2,h} \|u\|_{2,h}$$

From (29), we get

$$\|u\|_{1,h}^2 \leq \|f\|_{2,h} \|u\|_{2,h} \quad (30)$$

Stability

Lemma

There exists a constant positive C_Ω such that

$$\|u\|_{2,h} \leq C_\Omega \|u\|_{1,h}$$

Proof: Since $u_0 = 0$ then

$$u_i = \sum_{i'=1}^i (u_{i'} - u_{i'-1}) = \sum_{i'=1}^i \frac{u_{i'} - u_{i'-1}}{h} \cdot h = \sum_{i'=1}^i (D_x u)_{i'} \cdot h$$

Thus

$$u_i^2 \leq \sum_{i'=1}^i h \sum_{i'=1}^i (D_x u)_{i'}^2 h \leq \sum_{i'=1}^{N-1} (D_x u)_{i'}^2 h = \|u\|_{1,h}^2$$

Stability

So

$$\|u\|_{2,h}^2 = \sum_{i=1}^{N-1} h u_i^2 \leq \sum_{i=1}^{N-1} h \|u\|_{1,h}^2 = h(N-1) \|u\|_{1,h}^2 \leq \|u\|_{1,h}^2$$

We have completed the proof of the lemma. Using the lemma and (30), we get

$$\|u\|_{1,h} \leq \|f\|_{2,h}$$

Consistency

Let L be the differential operator, \hat{u} be a exact solution of the following equation:

$$Lu(x) = f(x), \text{ for all } x \in \Omega$$

Let L_h be the discrete differential operator of L , and u be the discrete solution, we have

$$L_h u_i = f_i \text{ for all } i \in [1, N - 1]$$

Consistency (Cont.)

Definition

A finite differential scheme is said to be consistent with the partial differential equation it present, if for any smooth solution u , the truncation error of the scheme:

$$\tau_i = L_h \hat{u}(x_i) - f(x_i) \text{ for all } i \in [1, N - 1]$$

tends uniformly forward to zero when h tends to zero, that mean that

$$\lim_{h \rightarrow 0} \|\tau\|_{\infty, h} = 0$$

Consistency (Cont.)

Lemma

Suppose $\hat{u} \in C^4(\Omega)$. Then, the numerical scheme in (10) is consistent and second-order accuracy for the norm $\|\cdot\|_\infty$

Proof: We write again the definition L , L_h operators of our case:

$$L(\hat{u})(x_i) = -\frac{\partial^2 \hat{u}}{\partial x^2}(x_i)$$
$$L_h(\hat{u})(x_i) = -\frac{\hat{u}(x_{i-1}) - 2\hat{u}(x_i) + \hat{u}(x_{i+1}))}{h^2}$$

By using the fact that

$$L(\hat{u})(x_i) = -\frac{\partial^2 \hat{u}}{\partial x^2}(x_i) = f(x_i)$$

Consistency (Cont.)

We have

$$\tau_i = L_h(\hat{u})(x_i) - f(x_i) = L_h(\hat{u})(x_i) - L(\hat{u})(x_i)$$

Using the definition of L and L_h , there holds

$$\tau_i = -\frac{\hat{u}(x_{i-1}) - 2\hat{u}(x_i) + \hat{u}(x_{i+1}))}{h^2} + \frac{\partial^2 \hat{u}}{\partial x^2}(x_i)$$

Using the Taylor series expansion respect x , there exists

$\eta_i \in [x_{i-1}, x_{i+1}]$ such that

$$-\frac{\hat{u}(x_{i-1}) - 2\hat{u}(x_i) + \hat{u}(x_{i+1}))}{h^2} + \frac{\partial^2 \hat{u}}{\partial x^2}(x_i) = \frac{-h^2}{12} \frac{\partial^4 \hat{u}}{\partial x^4}(\eta_i)$$

Consistency (Cont.)

we get

$$\tau_i = -\frac{h^2}{12} \frac{\partial^4 \hat{u}}{\partial x^4}(\eta_i) = -\frac{h^2}{12} \frac{\partial^2 f}{\partial x^2}(\eta_i)$$

Thus,

$$\|\tau\|_{\infty, h} \leq \frac{h^2}{12} \left\| \frac{\partial^2 f}{\partial x^2} \right\|_{\infty}$$

and

$$\|\tau\|_{2, h} \leq \frac{h^2}{12} \left\| \frac{\partial^2 f}{\partial x^2} \right\|_{2, h}$$

Convergence

Lemma

Let u be the exact solution and u_h be the discrete solution, there holds

$$\lim_{h \rightarrow 0} \|\hat{u} - u\|_{1,h} = 0.$$

Proof: We have

$$\tau_i = L_h(\hat{u})(x_i) - f(x_i) = L_h(\hat{u})(x_i) - L_h(u)(x_i) = L_h(\hat{u} - u)(x_i)$$

Using the proof of stability, we have

$$\|\hat{u} - u\|_{1,h} \leq \|\tau\|_{2,h} \leq \frac{h^2}{12} \left\| \frac{\partial^2 f}{\partial x^2} \right\|_{2,h}$$