



# Chapter 4:

## Time Varying Fields

## 4.1: Time Varying Fields and Potentials :

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## a) Introduction to Time Varying Fields

- ❖ Electric charges induce electric fields and electric current induce magnetic fields.
- ❖ the cases for *magnetostatic* and *electrostatic* where the magnetic field and the electric field are *constant with time*.
- ❖ In the *magnetostatic* and *electrostatic* cases, the  $\vec{E}$  and  $\vec{D}$  fields are *independent* of  $\vec{B}$  and  $\vec{H}$  fields.
- ❖ if the charge and current sources were to *vary with time t*,
  - not only will the fields also vary with time .
  - the electric and magnetic fields become interconnected .
- ❖ And the coupling between them produces electromagnetic waves capable of traveling through free space and in material media .

# ❖ Model :

## Maxwell's equations

$$\text{rot } \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t} \quad (1)$$

$$\text{rot } \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad (2)$$

$$\text{div } \vec{D} = \rho_V \quad (3)$$

$$\text{div } \vec{B} = 0 \quad (4)$$

$$\text{div } \vec{J} = -\partial \rho_V / \partial t \quad (5)$$

## Constitutive relations

$$\vec{B} = \mu \vec{H} = \mu_0 (\vec{H} + \vec{M})$$

$$\vec{D} = \epsilon \vec{E} = \epsilon_0 \vec{E} + \vec{P}$$

$$\vec{J} = \sigma \vec{E}$$

## Boundary conditions

$$H_{1t} - H_{2t} = J_S$$

$$E_{1t} - E_{2t} = 0$$

$$D_{1n} - D_{2n} = \rho_S$$

$$B_{1n} - B_{2n} = 0$$

$$J_{1n} - J_{2n} = -\frac{\partial \rho_S}{\partial t}$$

## b) Time Varying Potentials :

### 1. Magnetic vector potential :

$$\begin{cases} \operatorname{div} \vec{B} = 0 & (4) \\ \operatorname{div}(\operatorname{rot} \vec{A}) = 0 & \text{(vector algebra)} \end{cases}$$



$$\vec{B} = \operatorname{rot} \vec{A}$$

### 2. Electric scalar potential: (2) : $\operatorname{rot} \vec{E} = -\frac{\partial \vec{B}}{\partial t} = -\operatorname{rot} \frac{\partial \vec{A}}{\partial t}$

$$\begin{cases} \operatorname{rot}(\vec{E} + \frac{\partial \vec{A}}{\partial t}) = 0 \\ \operatorname{rot}(\operatorname{grad} \varphi) = 0 & \text{(vector algebra)} \end{cases}$$



$$\vec{E} = -\operatorname{grad} \varphi - \frac{\partial \vec{A}}{\partial t}$$

### 3. The Lorenz condition : multivalued $\rightarrow$ singlevalued

$$\operatorname{div} \vec{A} + \mu\epsilon \frac{\partial \varphi}{\partial t} = 0$$

### c) D'Alembert equation for vector potential:

$$\text{❖ (1) : } \vec{\text{rot}} \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t} \quad \Rightarrow \quad \vec{\text{rot}} \vec{B} = \mu \vec{J} + \mu \epsilon \frac{\partial \vec{E}}{\partial t}$$

$$\Rightarrow \vec{\text{rot}}(\vec{\text{rot}} \vec{A}) = \mu \vec{J} + \mu \epsilon \frac{\partial}{\partial t} (-\text{grad} \varphi - \frac{\partial \vec{A}}{\partial t})$$

$$\Rightarrow \text{grad}(\text{div} \vec{A}) - \Delta \vec{A} = \mu \vec{J} - \text{grad}(\mu \epsilon \frac{\partial \varphi}{\partial t}) - \mu \epsilon \frac{\partial^2 \vec{A}}{\partial t^2}$$

$$\text{❖ Using Lorenz condition : } \text{div} \vec{A} + \mu \epsilon \frac{\partial \varphi}{\partial t} = 0$$

➤ D'Alembert equation for vector potential:

$$\Delta \vec{A} - \mu \epsilon \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu \vec{J}$$

## d) D'Alembert equation for scalar potential:

❖ (3) :

$$\rho_V = \operatorname{div} \vec{D} = \varepsilon \cdot \operatorname{div} \left( -\operatorname{grad} \varphi - \frac{\partial \vec{A}}{\partial t} \right) = -\varepsilon \Delta \varphi - \varepsilon \frac{\partial}{\partial t} (\operatorname{div} \vec{A})$$

❖ Using Lorenz condition:  $\operatorname{div} \vec{A} + \mu \varepsilon \frac{\partial \varphi}{\partial t} = 0$

$$\Rightarrow \rho_V = -\varepsilon \Delta \varphi + \mu \varepsilon^2 \frac{\partial^2 \varphi}{\partial t^2}$$

➤ D'Alembert equation for scalar potential:

$$\Delta \varphi - \mu \varepsilon \frac{\partial^2 \varphi}{\partial t^2} = -\frac{\rho_V}{\varepsilon}$$

## ❖ Summary :

- i. Potentials  $\phi(t)$  and  $\vec{A}(t)$  satisfy wave equations .

$$\Delta \vec{A} - \frac{1}{v^2} \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu \vec{J}$$

$$\Delta \phi - \frac{1}{v^2} \frac{\partial^2 \phi}{\partial t^2} = -\frac{\rho_V}{\epsilon}$$

⇒ Time Varying Fields propagate in medium with  $v = \frac{1}{\sqrt{\mu\epsilon}}$

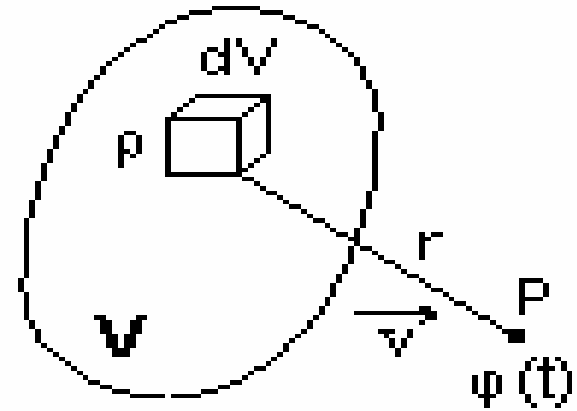
⇒ produce electromagnetic waves ⇒ Apply for communications



## ii. The solution of the wave equations :

$$\vec{A}(t) = \frac{\mu}{4\pi} \int_V \frac{\vec{J}(t - r/v) dV}{r}$$

$$\phi(t) = \frac{1}{4\pi\epsilon} \int_V \frac{\rho_V(t - r/v) dV}{r}$$



⇒  $\phi(t)$  and  $\vec{A}(t)$  : Retarded potentials